# SINGLE-VALUEDNESS OF A GENERAL SOLUTION OF THE PROBLEM OF MOTION OF A HEAVY RIGID BODY IN A NEWTONIAN FORCE FIELD IN ONE PARTICULAR CASE 

(OB ODNOZNACHNOSTI OBSHCHEGO RESHENIIA ZADACHI O dVIZHENII TIAZHELOGO TELA V N'IUTONOVSKOM POLE

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Iu. A. ARKHANGEL'SKII<br>(Moscow)<br>(Received September 3, 1964 )

It is well known that the equations of motion of a heavy rigid body about a fixed point in a Newtonian force field which are

$$
A \frac{d p}{d t}+(C-B) q r=-M g\left(y_{0} \gamma^{n}-z_{0} \gamma^{\prime}\right)+\frac{3 g}{R}(C-B) \gamma^{\prime} \gamma^{\prime \prime}, \quad \frac{d \gamma}{d t}=r \gamma^{\prime}-q \gamma^{2}
$$

can be reduced to quadratures in the real time only in two cases [ $1-3$ ]

$$
\begin{equation*}
x_{0}=0, \quad y_{0}=0, \quad z_{0}=0 ; \quad A=B, \quad x_{0}=0, \quad y_{0}=0 \tag{0.2}
\end{equation*}
$$

for which the fourth algebraic integral $[4,5]$ exists.
We have shown in [6] that all the possible cases of single-valuedness of the general solution of ( 0.1 ), are equivalent to the two cases shown in ( 0.2 ). From this it follows that the proof of the single-valuedness of the general proof for the above two cases (0.2), will clarify the question of the possibility of existence of the following general theorem [7]:

The fourth algebraic integral of the system (0.1) exists if, and only if general solations for $p, q, r, \gamma, \gamma^{\prime}, \gamma^{\prime \prime}$ exist, which are single-valned over the whole of the $t$-plane.

Proof of the single-valuedness of the general solution of (0.1) and of the remaining aix direction cosines in the second of the ( 0.2 ) cases, follows.

1. The case under consideration analogons to the Lagrange case in the classical problem of the motion of a heavy, rigid body about a fixed point, has the following equations of motion:

$$
\begin{gather*}
d p / d t-m q r=-(\alpha+a u) \gamma^{\prime}, \quad u=\gamma^{\prime \prime} \\
d q / d t+m p r=(\alpha+a u) \gamma, \quad d r / d t=0  \tag{1.1}\\
d \gamma / d t=r \gamma^{\prime}-q u \quad\left(p q r, \gamma \gamma^{\prime} u\right)
\end{gather*}
$$

four independent first integrals of which are

$$
\begin{gather*}
p^{2} \not q^{2}-2 \alpha u-a u^{2}=h, \quad p r+q \gamma^{\prime}-b u=k, \quad r=r_{0}, \quad \gamma^{2} \downarrow \gamma^{\prime 2}+u^{2}=1 \\
\alpha=-M g z_{0} / A, \quad m=(A-C) / A, \quad a=3 g m / R, \quad b=(m-1) r_{0} \tag{1.2}
\end{gather*}
$$

Introducing polar coordinates $\rho$ and $\sigma$

$$
\begin{equation*}
p=\rho \cos \sigma, \quad q=\rho \sin \sigma, \quad \rho^{2}=h+2 \alpha u \not+a u^{2} \tag{1.3}
\end{equation*}
$$

we obtain

$$
\begin{gather*}
(d u / d t)^{2}=\left(1-u^{2}\right) \rho^{2}-\overline{(k+b u)^{2} \equiv \Phi(u)} \\
d \sigma / d t=-r_{0}+\left(L_{1} u+L_{2}\right) / \rho^{2}, L_{1}=a k-a b, \quad L_{2}=\alpha k-b h \tag{1.4}
\end{gather*}
$$

from which $u$ and $\sigma[3]$ can be determined.
From the first equation of (1.4) it follows that $u$ can be found by transforming an elliptic integral and that it can, together with its derivative, be expressed in terms of singlevalued functions of time.

Taking into account (1.3) let $u s$, instead of $p$ and $q$, consider the function $W=\rho e^{i \sigma}$.
In order to show that $W$ is a single-valued function of time (from which the singlevaluedness of $p$ and $q$ also follows), we shall now prove that the function $W_{1}=e^{i r_{0} t} W$ is a single-valued function of $u$ and of $d u / d t$, and we shall do it by investigating its behavior near its singular points. We shall represent $W_{1}$ by (1.4) and

$$
\begin{equation*}
W_{1}=\rho e^{i \sigma_{1}} ; \quad \sigma_{1}=\int F_{1}(u) d u, \quad F_{1}(u)=\frac{L_{1} u+L_{2}}{\rho^{2} \sqrt{\Phi(u)}} \tag{1.5}
\end{equation*}
$$

factorising the polynomials $\rho^{2}$ and $\Phi(u)$ thus
$\boldsymbol{\rho}^{2}=a\left(u-\lambda_{1}\right)\left(u-\lambda_{2}\right), \quad \Phi(u)=-a\left(u-\mu_{1}\right)\left(u-\mu_{2}\right)\left(u-\mu_{3}\right)\left(u-\mu_{4}\right)$
Since the single-valuedness of the solution when all the roots of $\Phi(u)$ are distinct implies the single-valuedness of the general solution of (1.1), we shall continue to assume that $\mu_{l}(l=1,2,3,4)$ are all distinct.
2. The relationships between the roots of the polynomials $\lambda_{j}, \mu_{l}(j=1,2 ; l=1,2,3,4)$ the values $u=\lambda_{j}, u=\mu_{l}$ of which are singular points of $W_{l}$, admit of the five following cases:
$1^{\circ}$. Let $\lambda_{1} \neq \lambda_{2}, \lambda_{j} \neq \mu_{l}$. In this case (1.4) yields

$$
\begin{array}{rll}
\left(L_{1} u+L_{2}\right) \rho^{-2}= & \sum_{j=1}^{j=2} 1 / 2\left(k+b \lambda_{j}\right)\left(u-\lambda_{j}\right)^{-1}, & \Phi^{-1 / 2}\left(\lambda_{j}\right)=i \varepsilon\left(k+b \lambda_{j}\right)^{-1}  \tag{2.1}\\
k+b \lambda_{j} \neq 0, & \operatorname{res}\left[F_{1}(u)\right]_{u=\lambda_{j}}=1 / 2 i \varepsilon & (\varepsilon= \pm 1, j=1,2)
\end{array}
$$

From (1.5), (1.6) and (2.1) it follows that the points $u=\lambda_{j}$ are removable singular points, hence in their vicinity, $W_{1}$ is a single-valued function of $u$. Near each of the points $\mu_{l}(l=1,2,3,4)$ the first equation of (1.4) will give us $F_{1}(u)=\xi_{l}^{-1 / 2} f_{1}\left(\xi_{l}\right) ; \quad \sigma_{1}=\xi_{l}^{1 / 2 f_{2}}\left(\xi_{l}\right)=\Phi^{1 / 2}(u) f_{3}\left(\xi_{l}\right)=f_{3}\left(\xi_{l}\right)(d u / d t)_{1}^{\eta}\left(\xi_{l}=u-\mu_{l}\right)$ where $f_{s}(\eta)(s=1,2, \ldots)$ are holomorphic functions of $\eta$.

From (1.5) and (2.2) it follows that in vicinity of the points $u=\mu_{l}, W_{1}$ is a singlevalued function of $u$ and $d u / d t$. Since in the remaining cases the proof is almost identical with the one just given, we shall give it only when it differs from the one already given.
$2^{0}$. Let $\lambda_{1}=\lambda_{2}=\lambda, \lambda \neq \mu_{l}$. At $u=\lambda$, we have

$$
\begin{gather*}
\left(L_{1} u \not L_{2}\right) \rho^{-2}=(k+\lambda b)(u-\lambda)^{-1}, \quad \Phi^{-1 / 2}(\lambda)=i \varepsilon(k \ngtr \lambda)^{-1}  \tag{2.3}\\
k+\lambda b \neq 0, \quad \rho^{2}=a(u-\lambda)^{2}, \quad \operatorname{res}\left[F_{1}(u)\right]_{u=\lambda}=i \varepsilon, \quad \varepsilon= \pm 1 \tag{2.4}
\end{gather*}
$$

$3^{\circ}$. Let $\lambda_{1} \neq \lambda_{2}, \lambda_{1}=\mu_{1}, \lambda_{2} \neq \mu_{8}(s=2,3,4)$. Here, by the first equation of (1.4), we have $b^{2}+k^{2}: \neq 0$.

Near the point $u=\mu_{2}$, we shall have
$F_{1}(u)=\xi_{1}{ }^{-1 / 2} f_{4}\left(\xi_{1}\right) ; \quad \sigma_{1}=\xi_{1}^{-1 / 2} f_{5}\left(\xi_{1}\right)=\Phi^{-1 / 2}(u) f_{6}\left(\xi_{1}\right)=f_{6}\left(\xi_{1}\right)(d u / d t)^{-1}(2.5)$
and $\quad \rho=\xi_{1}^{1 / 2} f_{7}\left(\xi_{1}\right)=\Phi^{1 / 2}(u) f_{8}\left(\xi_{1}\right)=(d u / d l) f_{8}\left(\xi_{1}\right) \quad\left(\xi_{1}=u-\mu_{1}\right)$
40. Let $\lambda_{1} \neq \lambda_{2}, \lambda_{1}=\mu_{1}, \lambda_{2}=\mu_{2}$. In this case $b=k=0$, and from the first equation of (1.4) together with (1.5) it follows that $\sigma_{1}=$ const, while $\rho$ near the points $u=\lambda_{j}$ ( $j=1,2$ ) will assume the form (2.6).
50. Let $\lambda_{1}=\lambda_{2}=\lambda=\mu_{1}$. By (2.3) the expansion near the point $u=\lambda$, will be of the form (2.5).

Hence, $p, q, r, y^{*}=u$ are single-valued functions of time, and by (1.1), $\gamma$ and $y^{\prime}$ have the same property.
3. We shall now show that the remaining six direction cosines are also single-valued functions of time. Using Eulerian angles $\theta, \varphi, \psi$ in the usual manner, it will be sufficient to show [8] that
$\cos \varphi \cos \psi, \quad \cos \varphi \sin \psi, \sin \varphi \sin \psi, \sin \varphi \cos \psi, \sin \psi \sin \theta, \cos \psi \sin \theta$
are single-valued. From $\gamma, \gamma^{\prime}, \gamma^{\prime \prime}$

$$
r=\sin \theta \sin \varphi, \quad \gamma^{\prime}=\sin \theta \cos \varphi, \quad \gamma^{\prime \prime}=\cos \theta
$$

which were shown above to be single-valued functions of time,

$$
\begin{equation*}
\sin \varphi=V_{1} \sin \theta, \quad \cos \varphi=V_{2} \sin \theta \tag{3.2}
\end{equation*}
$$

follows, where $V_{s}(s=1,2, \ldots)$ are single-valued functions of time.
Instead of $\sin \theta \sin \psi$ and $\sin \theta \cos \psi$, let us consider

$$
\begin{equation*}
W_{2}=\sqrt{1-u^{2} e^{i} \psi} \tag{3.3}
\end{equation*}
$$

where $\psi$ by (1.2), can be determined from

$$
\begin{equation*}
d \psi / d t=(k+b u) /\left(1-u^{2}\right) \tag{3.4}
\end{equation*}
$$

We shall show that $W_{2}$ is a singlo-valued function of time. To do this, it will be sufficient to show that $W_{2}$ is a single-valued function of $u$ and $d u / d t$, where $\psi$, by the first equation of (1.4) and (3.4) is given by

$$
\begin{equation*}
\psi=\int F_{2}(u) d u, \quad F_{2}(u)=\frac{k+b u}{\left(1-u^{2}\right) \sqrt{\Phi(u)}} \tag{3.5}
\end{equation*}
$$

The proof differs from that given for $W_{1}$, in the following details:- in the formulas of the $1-s t, 3$-rd and 4 -th case, $\lambda_{1}$ and $\lambda_{1}$ should be replaced by 1 and -1 respectively, $\sigma_{2}$ replaced by $\psi, \rho$ by $\rho_{1}=\left(1-u^{2}\right)^{1 / 2}, F_{1}(u)$ by $F_{2}(u), L_{1}$ by $b$ and $L_{1}$ by $k$.

Hence $W_{2}$ is a single-valued function of time, and

$$
\begin{equation*}
\sin \psi=V_{3} \sin \theta, \quad \cos \psi=V_{4} \sin \theta \tag{3.6}
\end{equation*}
$$

From (3.2) and (3.6) it follows that (3.1) and hence the remaining six direction cosines, are single-valued functions of time.

It follows that in order to complete the existence proof of the theorem given above, the single-valuedness of the general solution in the first case of ( 0.2 ) with the condition $A \neq B$, remains to be investigated (for the case $A=B$, single-valuedness of the general solution follows from the results obtained above).

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